17 — Line Profiles II [Revision : 1.2]

- The Voigt function
 - In most general case, shape of line profiles is governed by a combination of natural/pressure/collisional broadening and Doppler broadening:

$$\phi(\nu) = \sqrt{\frac{\mu}{2\pi kT}} \int_{-\infty}^{\infty} e^{-\mu v^2/2kT} \frac{\gamma/4\pi^2}{(\nu + \nu v/c - \nu_0)^2 + (\gamma/4\pi)^2} dv$$

- Looks messy; but let

$$u = \frac{(\nu - \nu_0)}{\Delta \nu_{\rm D}}$$

(measure of distance from line center, in units of Doppler FWHM),

$$y = v \sqrt{\frac{\mu}{2kT}}$$

(measure of particle speed),

$$a = \frac{\gamma}{4\pi\Delta\nu_{\rm D}}$$

- (measure of Lorenz width vs. Doppler width)
- Then, profile is

$$\phi(\nu) = \frac{\sqrt{\pi}}{\Delta\nu_{\rm D}} H(a, u)$$

where

$$H(a, u) = \frac{a}{\pi} \int_{-\infty}^{\infty} \frac{e^{-y^2}}{(u-y)^2 + a^2} \, \mathrm{d}y$$

is the Voigt function \mathbf{V}

- Voigt function looks like Gaussian for small u (i.e., near line center), and like Lorentzian for large u (i.e., out in line wings)
- Schuster-Schwarzschild model
 - Simple model for the formation of spectral lines
 - Assumes lines are formed by cold reversing layer sitting above hot background source of radiation
 - In reversing layer, $S_{\lambda} \to 0$, and so emergent intensity is

$$I_{\lambda} = I_0 \mathrm{e}^{-\tau_{\lambda}}$$

where I_0 is background continuum intensity, and τ_{λ} is optical thickness of reversing layer

– Assume reversing layer is uniform with thickness Δ_z ; then

$$\tau_{\lambda} = \kappa_{\lambda} \rho \Delta_z = n \sigma(\nu) \Delta_z = \sigma(\nu) N$$

where $N = n\Delta_z$ is **column density** — number of particles per unit area that participate in the bound-free absorption responsible for a given spectral line

- Using general expression

$$\sigma(\nu) = \frac{e^2}{4\epsilon_0 mc} f \phi(\nu)$$

with $\phi(\nu)$ given by Voigt-function expression, result is

$$I_{\lambda} = I_0 \exp\left[\frac{e^2}{4\epsilon_0 mc} f N \frac{\sqrt{\pi}}{\Delta \nu_{\rm D}} H(a, u)\right]$$

- General behavior
 - * When N is small, damping coefficient γ is also small (i.e., little collisional/pressure broadening), and so profile looks Gaussian
 - * As N gets bigger, the optical depth at line center becomes so large that $I_{\lambda} = 0$ there (saturation)
 - * As N continues to get bigger, amount of collisional/pressure broadening increases, and so Lorentzian-dominated **damping winds** appear
- The Curve of Growth
 - From Schuster-Schwarzschild model, the equivalent width W of a profile shows three distinct regimes:
 - * $W \propto N$ for small column densities (profile not yet saturated)
 - * $W \propto \sqrt{\ln N}$ for intermediate column densities (profile saturated)
 - * $W \propto \sqrt{N}$ for large column densities (damping wings appear)
 - A plot of $\log_{10} W$ against $\log_{10} N$ shows these three regimes as the **curve of growth** (i.e., how the strength of a line grows as the column density increases)
 - Curve-of-growth analysis can be used to determine number of absorbing atoms of given type in atmosphere — useful for determining temperatures or abundances